

The Crystal Structure of Potassium Palmitate (Form B)

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Form B of potassium palmitate, $C_{16}H_{31}O_2K$, is triclinic with $a=4.15$, $b=5.60$, $c=37.82$ Å; $\alpha=93.0$, $\beta=91.4$, $\gamma=92.4^\circ$. The space group is PI with two molecules per unit cell. The hydrocarbon chains are all parallel to each other and are packed in accord with a somewhat distorted version of the usual triclinic sub cell. The ionic layers have essentially the same structure as those in form A of potassium caprate. The angle of tilt of the chain axes is 52° .

Introduction

Early work on the crystallography of potassium soaps was carried out by Piper (1929), who reported one form each for neutral and 1:1 acid potassium soaps. Later work by Vand, Lomer & Lang (1947) established the existence of three anhydrous forms A, B and C for neutral even-numbered potassium soaps, the B form being identical with that described by Piper. The unit-cell dimensions of potassium soaps in these three forms have been determined from powder photographs by Lomer (1952). Form A is monoclinic and occurs when soaps with from six to ten carbon atoms are crystallized from alcohol at room temperature; form B, which is triclinic, occurs in similar conditions for soaps of greater chain length. Both the A and B forms transform to the monoclinic C form at higher temperatures. The crystal structure of potassium caprate, form A, has been described by Vand, Lomer & Lang (1949). The determination of the structure of potassium palmitate, form B, is described below.

Experimental

Samples of potassium laurate, myristate, palmitate and stearate were prepared by neutralizing alcoholic solutions of the appropriate fatty acids with a solution of potassium hydroxide in alcohol which contained about 5% water. Phenolphthalein was used as an indicator. The solutions were then allowed to evaporate slowly in an evacuated desiccator containing calcium chloride. The solutions of the palmitate alone produced single crystals suitable for X-ray examination: as usual with these long-chain compounds the crystals were extremely thin plates, the large faces being (001). A crystal approximately $1\text{ mm} \times 0.5\text{ mm} \times 0.03\text{ mm}$ was mounted in a Weissenberg camera and rotation and moving film photographs were taken with the crystal rotating successively about [100] and [010]. The unit cell was found to be triclinic with $a=4.15 \pm 0.01$, $b=5.60 \pm 0.02$, $c=37.82 \pm 0.04$ Å, $\alpha=93.0 \pm 0.3$, $\beta=91.4 \pm 0.3$, $\gamma=92.4 \pm$

0.3° ; the long spacing ($=1/c^*$) was 37.80 Å, and assuming two molecules of $C_{16}H_{31}O_2K$ per unit cell the calculated density was $1.109 \pm 0.008\text{ g.cm}^{-3}$, which agrees satisfactorily with the observed value of $1.117 \pm 0.01\text{ g.cm}^{-3}$. The cell dimensions agree well with those determined earlier by Lomer (1952).

The intensities of the $h0l$ and $0kl$ reflexions were estimated visually: the multiple film technique was used in conjunction with a standard set of spots produced by timed exposures of one of the reflections from the crystal. Some difficulty was experienced in the intensity estimation because of the wide variation in spot shape; the average discrepancy between the estimates of the two authors was equivalent to an R index of about 10%. Cochran's (1948) method was used to calculate the Lorentz and polarization factors. The observed structure factors were subsequently brought to the absolute scale by comparison with calculated values. No absorption corrections were applied, since, apart from a very few reflexions which occur for nearly grazing incidence or emergence, the path length of the beam in the crystal does not greatly exceed the thickness of the crystal and thus for the great majority of reflections μt is less than 0.2.

Determination of the structure

The short cell dimensions of the A and B forms are almost exactly the same except that the a axis is twice as long in the A form as it is in the B form. This suggested that the crossing of the chain axis direction in alternate layers, perpendicular to the a axis, which occurs in the A form was absent in the B form. It was thus possible to visualize a trial structure for the B form in which the packing in the ionic double layer was the same as that in the A form, and in which all the hydrocarbon chains were arranged as those in one of the layers of the A form. This structure also explained the great observed intensity of the reflexions 014 and 104, and was consistent with the Patterson projections along [100] and [010]. The trial structure was refined in the initial stages by the usual method of electron density projections, and in the later stages by use of Curtis's (1959) least-squares program. The

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scattering factor data of Berghuis, Haanappel, Poters, Loopstra, MacGillavry & Veenendaal (1955) were used for the heavy atoms and those of McWeeny (1951) for hydrogen atoms. After several cycles of refinement R was 15.5% for the $h0l$ reflexions and 17.6% for the $0kl$ reflexions. At this stage the positional parameters of all the atoms, except the hydrogens, had been reliably determined but difficulty was experienced in deciding the best isotropic thermal motion parameters for the carbon and oxygen atoms. Values of B ($= 8\pi^2\mu_s^2$) of 2, 3 or 4 Å² all gave roughly equal values for R . It was therefore decided to refine the three-dimensional data collected from the upper layer line photographs.

A separate scale factor for each layer was used to bring the observed structure factors to the absolute scale: these factors were revised once during the several cycles of least-squares refinement that were carried out. The program of Hodgson & Rollett (1963) was used and the computing was done on the Oxford University Mercury Computer. Towards the end of the refinement the hydrogen atoms, placed in calculated positions, were included in the calculation of the structure factors but their parameters were not refined. Isotropic temperature factors for each of the other atoms were used and refined simultaneously with the position parameters. The final R value was 11.6% and the total number of structure factors used was about 850. The parameters resulting from the last cycle of refinement are given in Table 1. The accuracy of these parameters may be judged from the standard deviations given in Table 2. The values of the observed and calculated structure factors are given in Table 3.

Discussion of the results

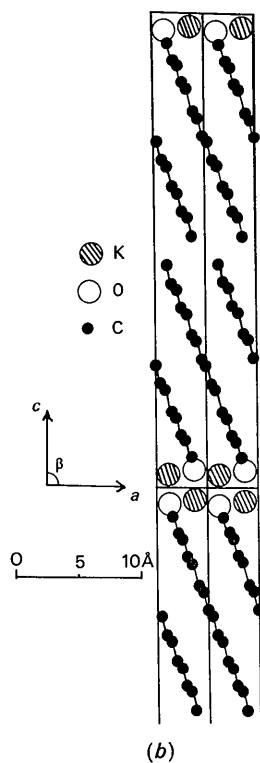
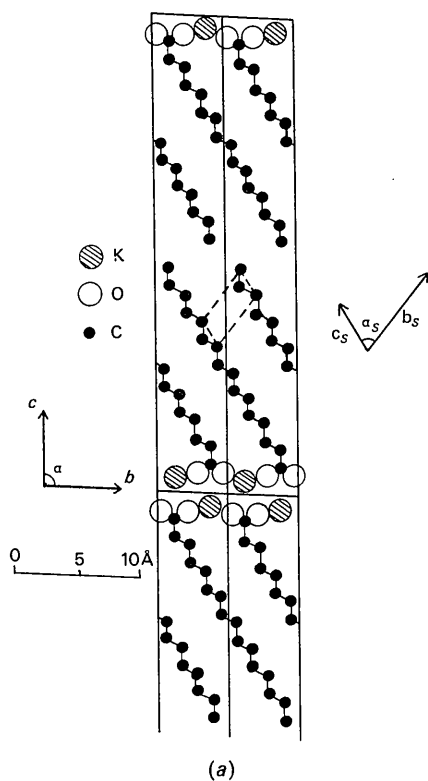
The general arrangement of the atoms in the crystal is illustrated in Fig. 1, which shows that, as in other soap structures, the potassium and oxygen atoms form double layers, parallel to (001), which are separated from each other by the hydrocarbon chains. The packing of the potassium and oxygen atoms in one half of a double layer is illustrated in Fig. 2; it differs from that found in the A form of potassium caprate chiefly in that the angle γ is 92.4° instead of 90°. Each potassium atom is surrounded by four oxygen atoms in the same half of the double layer and by two more in the other half of the double layer. The distances from a potassium atom to its four neighbours in the same half of the double layer are 2.71, 2.73, 2.73 and 2.78 Å and those to the two neighbours in the other half of the double layer are both 2.82 Å.

The carbon—oxygen bond lengths in the carboxyl group are respectively 1.21 ± 0.025 Å and 1.27 ± 0.025 Å and make an angle of 123° with each other. Both of these lengths are thus significantly shorter than the length of either a single bond or of a bond of partial double character and it is not clear why this should be so. The bond lengths between the carbon atoms

Table 1. Atomic coordinates as fractions of the unit cells

Atom	x/a	y/b	z/c	B
K	0.23761	0.25687	0.02755	1.96
O(1)	0.76645	0.57339	0.04122	3.1
O(2)	0.72621	0.96148	0.04114	3.4
C(1)	0.71214	0.77592	0.05632	2.8
C(2)	0.61077	0.77670	0.09534	2.9
C(3)	0.50343	0.54146	0.11026	3.1
C(4)	0.43396	0.56098	0.14913	3.0
C(5)	0.31384	0.32458	0.16358	3.0
C(6)	0.23401	0.33782	0.20269	3.6
C(7)	0.12604	0.10154	0.21710	3.3
C(8)	0.03194	0.11470	0.25574	3.2
C(9)	0.93194	0.88061	0.27044	3.3
C(10)	0.84033	0.89108	0.30903	3.6
C(11)	0.73586	0.65645	0.32386	3.5
C(12)	0.64259	0.67219	0.36263	4.1
C(13)	0.53943	0.43008	0.37654	4.0
C(14)	0.44560	0.44556	0.41574	4.3
C(15)	0.34819	0.20426	0.42890	5.8
C(16)	0.25107	0.21069	0.46828	7.2
H(1)	0.837	0.827	0.1103	5
H(2)	0.430	0.911	0.0983	5
H(3)	0.287	0.496	0.0956	5
H(4)	0.694	0.412	0.1076	5
H(5)	0.661	0.611	0.1640	5
H(6)	0.254	0.695	0.1520	5
H(7)	0.097	0.272	0.1488	5
H(8)	0.504	0.188	0.1608	5
H(9)	0.459	0.389	0.2178	5
H(10)	0.051	0.473	0.2058	5
H(11)	0.910	0.051	0.2022	5
H(12)	0.317	0.967	0.2142	5
H(13)	0.254	0.167	0.2709	5
H(14)	0.847	0.251	0.2589	5
H(15)	0.712	0.830	0.2560	5
H(16)	0.110	0.746	0.2680	5
H(17)	0.064	0.943	0.3244	5
H(18)	0.657	0.027	0.3124	5
H(19)	0.510	0.608	0.3092	5
H(20)	0.917	0.524	0.3212	5
H(21)	0.871	0.721	0.3779	5
H(22)	0.465	0.805	0.3659	5
H(23)	0.315	0.377	0.3618	5
H(24)	0.722	0.293	0.3738	5
H(25)	0.676	0.495	0.4314	5
H(26)	0.269	0.578	0.4194	5
H(27)	0.127	0.152	0.4140	5
H(28)	0.534	0.068	0.4260	5

were calculated and the average is 1.521 ± 0.010 Å; the average distance between alternate carbon atoms is 2.557 ± 0.005 Å and the C—C—C bond angle is $114.4 \pm 1^\circ$. The carbon atoms are all co-planar within the limits of experimental error, but the oxygen atoms lie one on either side of this plane. The displacements may be regarded as resulting from a 16.4° rotation of the carboxyl group about the first carbon—carbon bond in the molecule. A similar rotation of the carboxyl group has been reported in the crystal structure of the A₁ form of lauric acid by Lomer (1963) and also in the crystal structure of tetradecanamide by Turner & Lingafelter (1955) and in 11-aminoundecanoic acid hydrobromide hemihydrate by Sim (1955). These results suggest that this rotation may be a property of the fatty acid radical rather than due to an effect of the end packing of the molecules, which is different, both physically and ge-



ometrically in the ionically bonded soaps and the hydrogen bonded acid and amides. The angle of tilt (*i.e.* the angle between the chain axes and the (001) plane) is $52.03 \pm 0.20^\circ$. The individual bond lengths and angles, together with their calculated standard deviations are given in Fig. 4.

The side by side packing of the hydrocarbon chains can be described by a *subcell* (Vand & Bell, 1951; von Sydow, 1956) which has dimensions $a_s = 4.15$, $b_s = 5.30$,

Table 2. Standard deviations of the atomic coordinates

Atom	σ_x Å	σ_y Å	σ_z Å
K	0.0037 Å	0.0041 Å	0.0036 Å
O(1)	0.0104	0.0113	0.0102
O(2)	0.0107	0.0118	0.0106
C(1)	0.0157	0.0180	0.0156
C(2)	0.0157	0.0177	0.0156
C(3)	0.0160	0.0177	0.0159
C(4)	0.0160	0.0180	0.0159
C(5)	0.0160	0.0178	0.0158
C(6)	0.0170	0.0187	0.0167
C(7)	0.0167	0.0186	0.0162
C(8)	0.0164	0.0181	0.0160
C(9)	0.0166	0.0184	0.0161
C(10)	0.0170	0.0190	0.0167
C(11)	0.0172	0.0188	0.0166
C(12)	0.0179	0.0196	0.0177
C(13)	0.0177	0.0192	0.0173
C(14)	0.0185	0.0205	0.0182
C(15)	0.0212	0.0230	0.0209
C(16)	0.0255	0.0257	0.0240

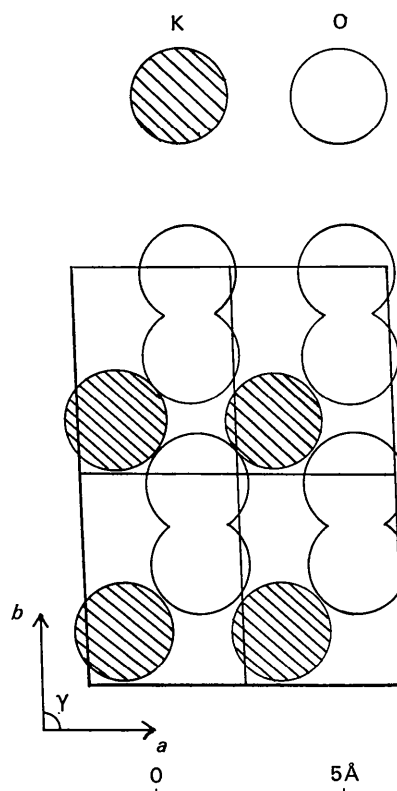


Fig. 1. Arrangement of the atoms in the crystal. (a) is the *a*-axis projection; the parallelogram in broken lines is the projection of the subcell; a_s is parallel to *a*. (b) is the *b*-axis projection of the structure.

Fig. 2. Arrangement of the atoms in one half of the ionic double layer.

Table 3. Values of observed and calculated structure factors

The structure factors are arranged in blocks of constant h and k; each block is headed with *hk. The first column gives the l index, the second column gives the observed value and the third column gives the calculated value. The values given are fifty times the absolute values.

Table with multiple columns containing numerical values for structure factors, organized in blocks by h and k indices. The table includes observed and calculated values for various l indices.

$c_s = 2.557 \text{ \AA}$; $\alpha_s = 65.2^\circ$, $\beta_s = 108.5^\circ$, $\gamma_s = 110.4^\circ$. The a_s axis of the subcell has the same orientation and length as the a axis of the crystal unit cell; the relation of the other subcell axes to those of the crystal cell is indicated in Fig. 1(a). These subcell dimensions differ considerably from those found in the hydrocarbon octadecane by Hayashida (1962). We attribute these differences to the effect of the strong binding in the ionic layers of the soap crystal, which must, to some extent, influence the packing of the chains. It would be interesting to measure the dimensions of the subcell in other triclinic soaps to see if this postulated influence of the ionic layers decreases as the length of the hydrocarbon chains increases. The packing of the methyl groups is illustrated in Fig. 3; each methyl group in the lower half of the cell is contact with four methyl groups in the upper half of the cell which are at distances from it of 4.35, 4.45, 4.01 and 4.02 \AA .

The thermal vibration parameters show that the atoms nearer to the ionic layer have smaller vibration amplitudes than those at the farther end of the molecules. This effect is not unexpected since the movement of these atoms is restrained by the heavy potassium atom. This observation explains, at least partly, why the atoms near the methyl group end of a hydrocarbon chain often appear, on electron-density maps, as peaks of lower height than those corresponding to atoms nearer the other end of the chain. We had hoped to investigate the anisotropy of the thermal motions in this compound but as the final residual of 11.5% was already as low as could be expected from the estimated errors of our intensity estimation, it did not seem

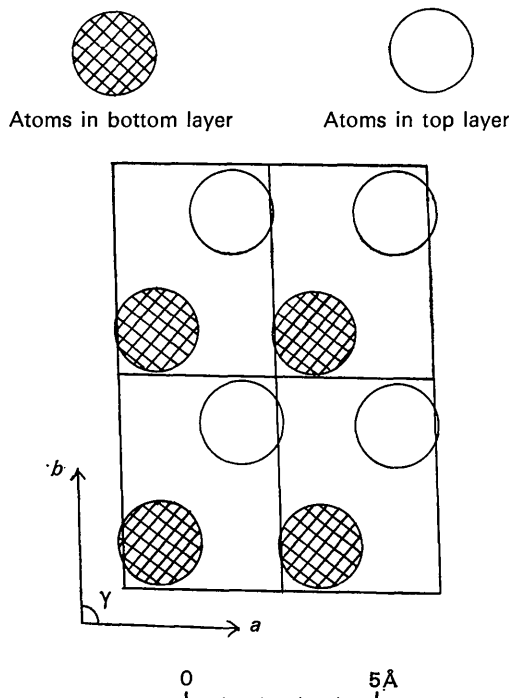


Fig. 3. Arrangement of methyl groups in the crystal.

reasonable to expect to obtain reliable values for any such anisotropy.

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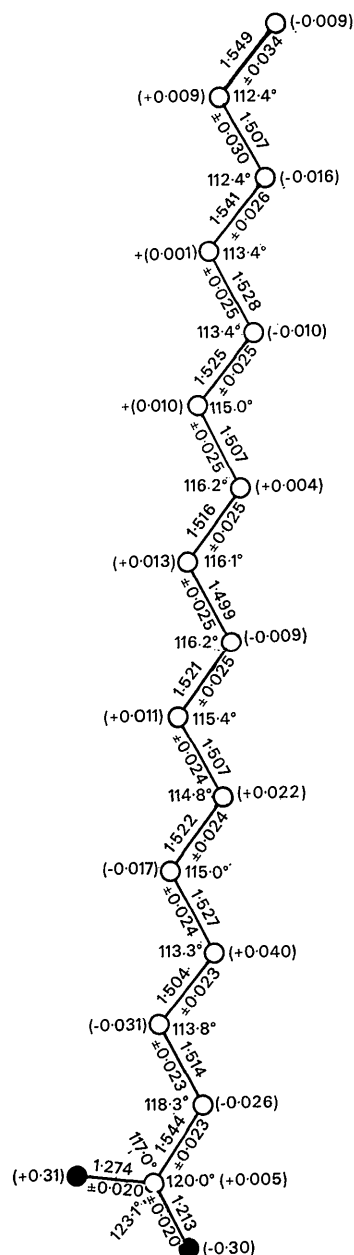


Fig. 4. Arrangement of atoms in the palmitate ion. Open circles are carbon atoms; full circles are oxygen atoms. The figures in parentheses are the deviations (\AA) of the atoms from the best plane through the carbon atoms. The bond lengths and their calculated standard deviations are given in \AA . The calculated standard deviations of the angles are all approximately 1.0° .

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A Refinement of the Crystal Structure of Yttria

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The crystal structure of yttria, Y_2O_3 , has been refined using single-crystal X-ray diffraction data. The results confirm the general features of earlier analyses using powder diffraction data, but a considerable improvement in accuracy has been achieved.

Introduction

Yttria, Y_2O_3 , is one of a series of metal oxides which have similar face-centered cubic structures conforming to the space group $Ia\bar{3} (T_h^7)^*$. Eight metal ions are in the special positions $\frac{1}{2}, \frac{1}{2}, \frac{1}{2}$; the remaining twenty-four occupy the sites $u, 0, \frac{1}{4}$. The forty-eight oxygen ions are in general positions x, y, z , and are arranged in distorted octahedra around the metal ions, the metal-oxygen bonding distances being unequal. This structure is generally referred to as the C-type metal oxide structure. All the sesquioxides of the rare earths belong to this system, as do Ga_2O_3 , In_2O_3 , Tl_2O_3 and the mineral bixbyite, $(Fe, Mn)_2O_3$.

The first basically correct determination of the structure of a C-type metal oxide was carried out by Pauling & Shappel (1932), using a trial-and-error method to analyse the powder-diffraction pattern of bixbyite. The accuracy of the analysis was low, for there are relatively few non-overlapping lines in the powder pattern. Moreover, the atomic scattering factor curves were quoted as having a possible error of twenty per cent. The oxygen positions were not determined directly, but from packing and bonding considerations they were deduced to be at sites equidistant from the metal ions. This structure was redetermined by Dachs (1956),

who made a least-squares analysis of the powder-diffraction pattern. The parameters obtained were quite different from those found by Pauling & Shappel.

The first experimental determination of the structure of yttria itself was reported by Villian (1957) who examined the neutron-diffraction powder pattern. A similar study was made by Fert (1962) who used a least-squares method on sixteen reflexions to find the parameters x, y, z, u as well as the neutron scattering length of yttrium.

Table 1 shows the results of these independent determinations, and it can be seen that there is a large range in the atomic coordinates quoted. Standard deviations on neither the atomic coordinates nor thermal parameters have been reported.

Table 1. *Structure parameters from previous investigations*

	u	x	y	z
Pauling & Shappel	0.030	0.385	0.145	0.380
Dachs	0.034	0.375	0.162	0.400
Villian	0.034	0.396	0.155	0.383
Fert	0.034	0.389	0.150	0.377

In view of the importance of the C-type metal oxide structure in oxide chemistry an analysis has been carried out with three-dimensional X-ray diffraction data collected from a single crystal. Since there is no overlapping of non-equivalent reflexions the number of independent observations is much greater than can be

* The space groups of several C-type crystals are reported incorrectly in the ASTM Card File as T_5 . The systematic absences of the reflexions for this space group are similar to those of T_h^7 except that the condition $hk0$ $h(k) = 2n + 1$ is allowed.